

Visibility of microwave absorption as a probe of Majorana bound states non-locality



Sarath Prem ^{1†}, Olesia Dmytruk ², Mircea Trif ¹

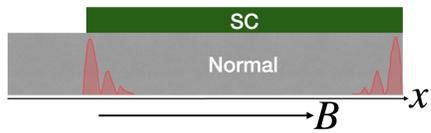


¹International Research Centre MagTop, Institute of Physics, Polish Academy of Sciences, Aleja Lotników 32/46, PL-02668 Warsaw, Poland

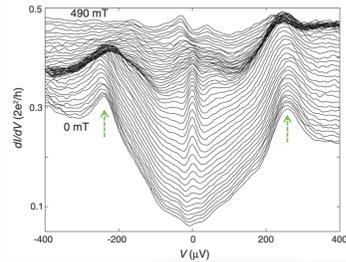
²CPHT, CNRS, École polytechnique, Institut Polytechnique de Paris, 91120 Palaiseau, France

†email: sarathprem@magtop.ifpan.edu.pl

Majorana bound state (MBS)



Topologically protected zero energy MBS (in red) in spin-orbit nanowire [1]



Zero-bias conductance peak: A predicted signature of Majorana bound states [2]

Ambiguity in the detection of topologically protected MBS:

- Trivial Andreev bound state (ABS) shows similar signature like zero-bias conductance peak and zero energy for a range of magnetic fields. [3]
- Quasi-Majorana bound state (QMBS) arise due to defects or imperfect gate voltages leading to "local topological states" which are difficult to control.

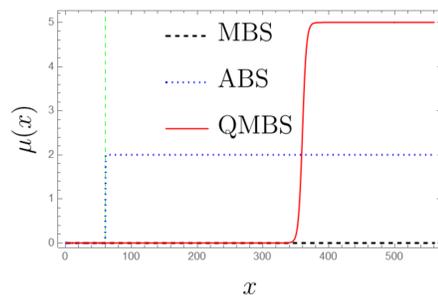
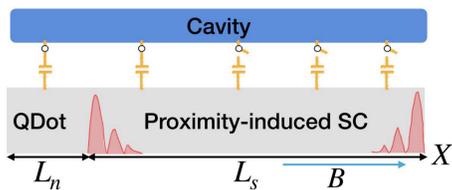
Challenge

Distinguish Majorana bound states from ABS [3, 4] and QMBS [4] as the latter two also give zero-bias conductance peak

Our solution

Probe the non-locality of Majorana bound state with microwaves!

Model Hamiltonian



(Left) A Rashba nanowire (gray part) with quantum dot (QDot) and partial proximity induced superconductivity is capacitively coupled to a cavity, and has a magnetic field \mathbf{B} aligned along the wire axis. (Right) Chemical potential profile with the green vertical line at QDot-SC interface.

- Model Hamiltonian in the tight-binding model is $H_T = H_{el} + H_c + H_{e-c}$ where,

$$H_{el} = \sum_{\sigma, \sigma'} \sum_{j=1}^L c_{j, \sigma}^\dagger [(2t_h(x_j) - \mu(x_j)) \delta_{\sigma\sigma'} + V_Z(x_j) \sigma_{\sigma\sigma'}^x] c_{j, \sigma'} + \sum_{j=1}^L (\Delta(x_j) c_{j, \uparrow}^\dagger c_{j, \downarrow}^\dagger + h.c.) - \sum_{\sigma, \sigma'} \sum_{j=1}^{L-1} [c_{j+1, \sigma}^\dagger (t_h(x_j) \delta_{\sigma\sigma'} - i\alpha(x_j) \sigma_{\sigma\sigma'}^y) c_{j, \sigma'} + h.c.],$$

$$H_{e-c} = \sum_{\sigma} \sum_j g(x_j) c_{j, \sigma}^\dagger c_{j, \sigma} (a^\dagger + a), \quad H_c = \omega_c a^\dagger a,$$

$c_{j, \sigma}, c_{j, \sigma}^\dagger$: electron operators with spin σ at site j , $\alpha(x_j)$: spin-orbit strength, $V_Z(x_j)$: Zeeman energy, $\mu(x_j)$: chemical potential, $\Delta(x_j)$: proximity superconducting gap, ω_c : cavity frequency, $a(a^\dagger)$: photon operators, and $g(x_j)$: cavity-wire coupling strength.

- We define **visibility of microwave absorption associated with the two parities** $\nu(\omega, x_j)$ [5] from left end of wire to position x_j of wire as

$$\nu(\omega, x_j) = \frac{\text{Im}[\chi_o(\omega, x_j)] - \text{Im}[\chi_e(\omega, x_j)]}{\text{Im}[\chi_o(\omega, x_j)] + \text{Im}[\chi_e(\omega, x_j)]},$$

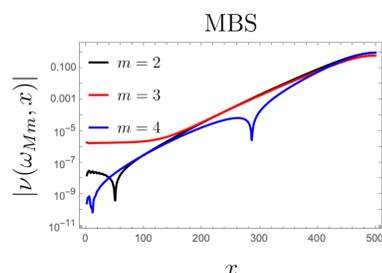
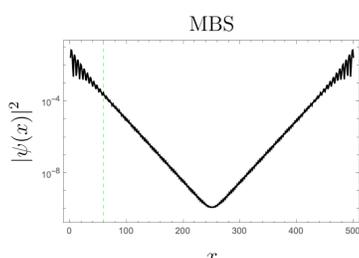
where $\chi_{o(e)}(\omega, x_j) = \sum_{\epsilon_m > 0} \left(\frac{|\mathcal{M}_{mM}^o(x_j)|^2}{\omega \pm \epsilon_m - \epsilon_m + i\eta} \right)$ is the susceptibility [6],

$\mathcal{M}_{mM}^e(x_j) = \sum_{\sigma, \rho=1}^{l=j} [u_{M\sigma}^*(x\rho) u_{m\sigma}(x\rho) - v_{M\sigma}^*(x\rho) v_{m\sigma}(x\rho)]$ and

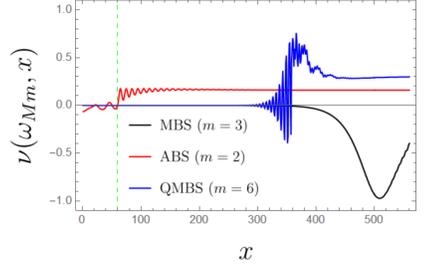
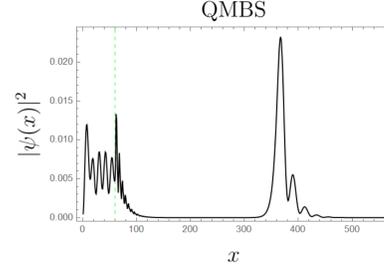
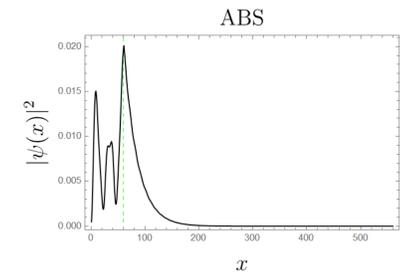
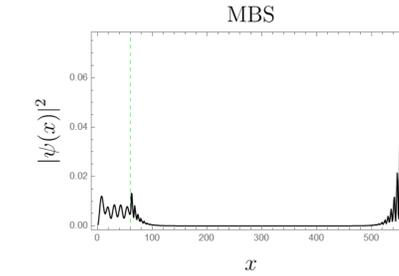
$\mathcal{M}_{mM}^o(x_j) = \sum_{\sigma, \rho=1}^{l=j} [v_{M\sigma}(x\rho) u_{m\sigma}(x\rho) - u_{M\sigma}(x\rho) v_{m\sigma}(x\rho)]$. Here u, v come from

$c_{j, \sigma}^\dagger = \sum_{n \geq 0} [u_{n\sigma}^*(x_j) b_n^\dagger + v_{n\sigma}(x_j) b_n]$ with b_n, b_n^\dagger as Bogoliubov operators.

- MBS in isolated SC has same exponential scaling behavior as visibility** $\nu(\omega, x_j)$:

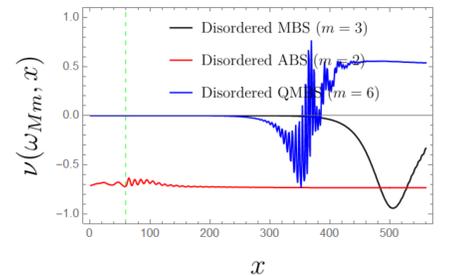
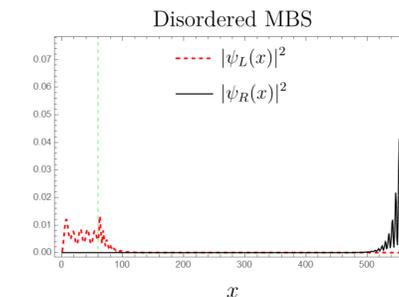


Majorana, Andreev, quasi-Majorana bound states



- MBS: Visibility captures the non-local nature of MBS as it peaks only when the cavity couples to both spatially separated wavefunction peaks.**
- ABS: In contrast to MBS, cavity couples to the peak that is localized near interface leading to visibility peak near interface.
- QMBS: Visibility peaks as the cavity couples to the peaks at edges of topological phase and not at SC edge.

Gaussian disorder in the SC part

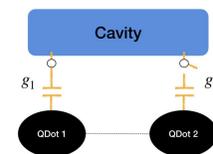


- Disorder strength less than minimal superconducting gap.
- MBS: Visibility retains its character due to the topological robustness of MBS against disorder.**
- ABS: Visibility is affected as ABS is not topologically protected against disorder.
- QBS: Visibility maintains its qualitative features.

Poor man's Majorana

- Poor man's Majorana [7]: Lack topological protection, but mimics Majorana properties

Visibility $\nu(\omega) = (2g_1 g_2) / (g_1^2 + g_2^2)$ where, $g_{1,2}$ is dot-cavity coupling strength, shows the non-local coupling nature as it requires $g_1 \neq 0$ and $g_2 \neq 0$.



Conclusions and outlook

- MBS has unique visibility profile due to its non-local nature that helps to distinguish it from ABS & QMBS.
- Visibility has same scaling behavior as MBS wavefunction (results not shown here).
- We also considered the effect of barrier (results not shown here).
- Possibility of applying this non-local visibility to other Majorana platforms?

References

- K. Laubscher and J. Klinovaja, J. Appl. Phys. **130**, 081101 (2021).
- V. Mourik et al., Science **336**, 1003 (2012).
- C. Reeg, O. Dmytruk, D. Chevallier, D. Loss, and J. Klinovaja, Phys. Rev. B **98**, 245407 (2018).
- E. Prada et al., Nat Rev Phys **2**, 575 (2020).
- P.-X. Shen, V. Perrin, M. Trif, and P. Simon, Phys. Rev. Research **5**, 033207 (2023).
- O. Dmytruk and M. Trif, Phys. Rev. B **107**, 115418 (2023).
- M. Leijnse and K. Flensberg, Phys. Rev. B **86**, 134528 (2012).



Acknowledgements: This research was partially supported by the "MagTop" project (FENG.02.01-IP.05-0028/23) carried out within the "International Research Agendas" program of the Foundation for Polish Science co-financed by the European Union under the European Funds for Smart Economy 2021-2027 (FENG).