

Known: 1d ultracold gases with $g>0$ host dark solitons in thermal equilibrium.

Question: under what conditions do they dominate the state?

$$\hbar \frac{d\Psi(x)}{dt} = (-i - \gamma) \mathcal{P}_{k_{\max}} \left[\left(-\frac{\hbar^2}{2m} \nabla^2 - \mu + g |\Psi(x)|^2 \right) \Psi(x) + \sqrt{2\gamma k_B T \hbar} \eta(x, t) \right]$$

PSGPE – projected stochastic Gross-Pitaevskii equation generates the thermal states

Two essential parameters of the uniform Bose gas:

Lieb-Liniger
interaction strength

$$\gamma_{LL} = \frac{mg}{\rho \hbar^2}$$

relative temperature
(T_d =degeneracy)

$$\tau = \frac{T}{T_d} = \frac{2k_B T m}{\hbar^2 \rho^2}$$

Scale to natural (healing length) units:

$$\xi_\rho = \hbar / \sqrt{mg\rho} = 1/(\rho \sqrt{\gamma_{LL}})$$

$$t_{\xi\rho} = \hbar/(g\rho) = 1/(\gamma_{LL}\rho^2)$$

m

And normalize to unit density at $T=0$:

$$\psi(x) = \frac{\phi(x)}{\gamma_{LL}^{1/4}}$$

Leads to: scaled PSGPE

$$\frac{d\phi(x)}{dt} = (-i - \gamma) \mathcal{P}_{k_{\max}} \left[\left(-\frac{1}{2} \nabla^2 - \mu + |\phi(x)|^2 \right) \phi(x) + \sqrt{\tau_{\text{dyn}}} \sqrt{\gamma} \eta(x, t) \right]$$

coupling to
high energy tails
generally small

chempot
sets density

thermal
complex noise

One essential dynamics parameter:
"dynamical" temperature

$$\tau_{\text{dyn}} = \frac{\tau}{\sqrt{\gamma_{LL}}} = \frac{2k_B T \sqrt{m}}{\hbar \rho^{3/2} \sqrt{g}}$$

Counting solitons, try 1: match dips to perfect soliton profile

Matching procedure:

- consider only lowest minima within a $\pm \xi$ range
- guess mean background density n_{est} , choose range r
- match actual minimum density $n_{0,\text{meas}}$ and depth d_{meas} to perfect soliton $n_{0,\text{perf}}, d_{\text{perf}}$ to get velocity u and background density $n\text{bar}$
- compare phase jump $\theta_{\text{pred}}(u, n\text{bar})$ predicted from u and $n\text{bar}$ to actual jump θ_{meas}
- count as soliton if e.g. $u < 0.5$ and the match is sufficient $0.75 < \theta_{\text{pred}}/\theta_{\text{meas}} < 1.25$
- compare mean accepted $n\text{bar}$ to the guess n_{est} , and redo with better n_{est} if needed.

$$d_{\text{pred}} = \bar{n}(1-u^2) \tanh^2 \left(r \sqrt{\frac{\bar{n}}{n_{\text{est}}}} \sqrt{1-u^2} \right)$$

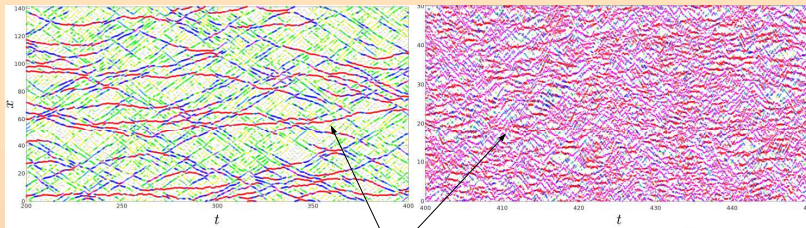
$$n_{0,\text{pred}} = \bar{n} u^2$$

$$\theta_{\text{pred}} = 2 \tan^{-1} \left[\frac{\sqrt{1-u^2}}{u} \tanh \left(r \sqrt{\frac{\bar{n}}{n_{\text{est}}}} \sqrt{1-u^2} \right) \right]$$

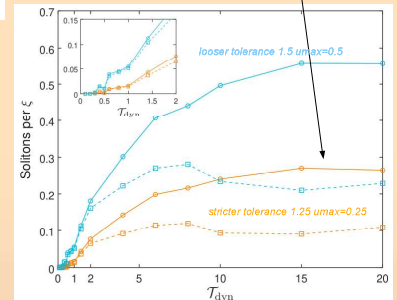
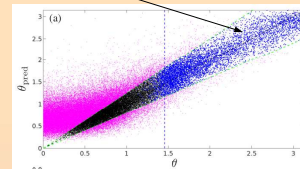
perfect gray soliton

Conclusions about counting:

- appearance of solitons at low T caught OK
- saturation by rubbish at high T
- need to consider soliton history



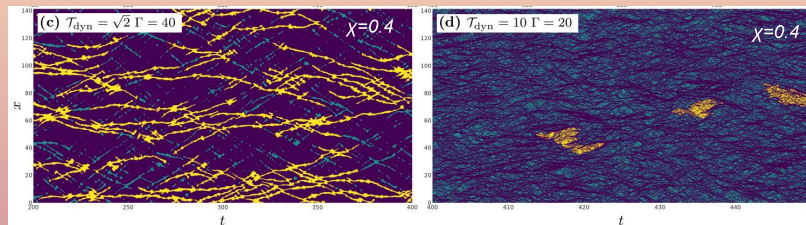
matches: red & blue: ok tolerance $0.75 < \theta_{\text{pred}}/\theta_{\text{meas}} < 1.25$ and speed $|u| < 0.25, 0.5$, green: fast, pink/cyan/yellow/gray: poor fits



Counting solitons, try 2: fill troughs, keep long-lived ones

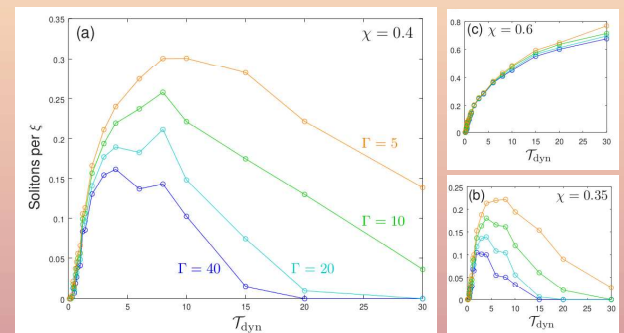
χ = max. density threshold "inside" soliton

Γ = minimum connected volume to treat as soliton



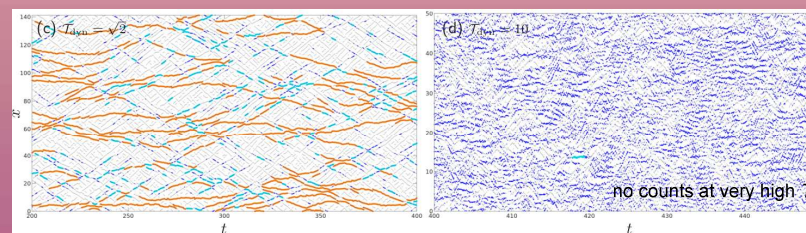
works well at low T

catches meaningless lumps at high T



problems again with characterising the high T end

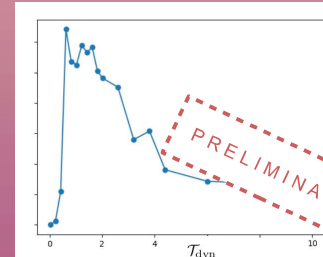
Counting solitons, try 3: keep only long clusters of dips



Clustering procedure:

- only consider dips deemed good enough (e.g. $0.67 < \theta_{\text{pred}}/\theta_{\text{meas}} < 1.5$ and speed $|u| < 0.5$, above)
- collect dips into chains, if dips at successive sample times are closer than $\epsilon \xi$ (above: $\epsilon=2/3$)
- Above: Orange – chains of length $> 5t_\xi$, cyan: length $= 2t_\xi$

e.g. soliton density in clusters of length $> 2t_\xi$



This approach appears to work as required.

- catches all solitons at low T
- ceases to count them at high T

$$0.3 \lesssim \tau_{\text{dyn}} = \frac{\tau}{\sqrt{\gamma_{LL}}} \lesssim \mathcal{O}(5)$$