

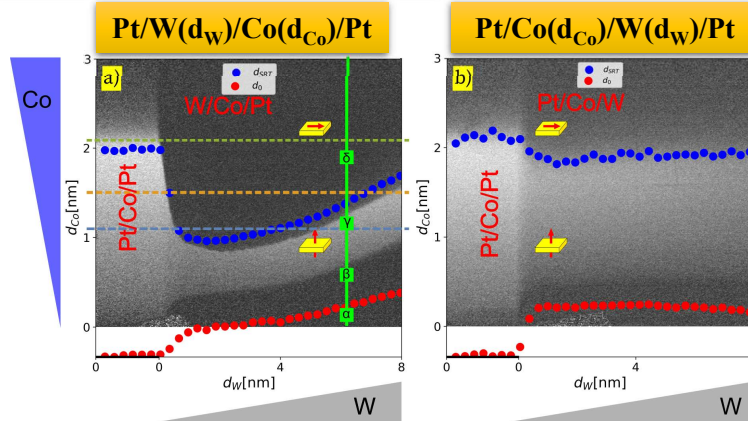
W. Dobrogowski<sup>1</sup>, R. Gieniusz<sup>1</sup>, U. Guzowska<sup>1</sup>, Z. Kurant<sup>1</sup>, I. Sveklo<sup>1</sup>, A. Pietruczik<sup>2</sup>, A. Maziewski<sup>1</sup>, A. Wawro<sup>2</sup>

<sup>1</sup>Department of Physics of Magnetism, Faculty of Physics, University of Białystok, 1L Ciołkowskiego St., 15-245 Białystok, Poland,

<sup>2</sup>Institute of Physics Polish Academy of Sciences, al. Lotników 32/46, 02-668 Warsaw, Poland

## Sample description

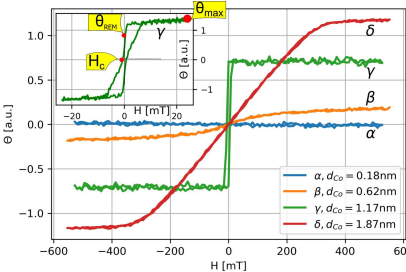
Ultra-thin ferromagnetic layers located between two different heavy metals are of great interest due to their potential application as a new material for magnetic memory or spin devices [1, 2]. The double wedge geometry samples Pt/W( $d_w$ )/Co( $d_{co}$ )/Pt and Pt/Co( $d_{co}$ )/W( $d_w$ )/Pt, where  $d_w$  and  $d_{co}$  are the thicknesses of W and Co layers, were epitaxially grown on sapphire substrate.



**Methodology** The influences of  $d_{co}$  and  $d_w$  on magnetization parameters were studied using polar magneto-optical Kerr effect (PMOKE) and Brillouin light scattering (BLS) methods. The dependence of PMOKE magnetization on applied in-plane field was used to determine uniaxial anisotropy field  $H_{A1}$ ;  $H_{A1} > 0$ ,  $H_{A1} < 0$  defining perpendicular and in-plane magnetization configurations, respectively.  $H_{A1}(d_{co}, d_w) = 0$  corresponds to spin reorientation transition (SRT) Co layer thicknesses. The dead layer thickness  $d_0$  was determined from the dependency of maximal PMOKE rotation  $\theta_{max}(d_{co})$  as interception of its linear approximation with horizontal axis (Fig. 3). The strength of interfacial Dzyaloshinskii-Moriya interaction (IDMI) and spin wave (SW) damping for selected  $d_{co}$  thicknesses as a function of  $d_w$  were determined from the dynamic characteristics measured by BLS in Damon-Eshbach mode. The inclination of the dependency of frequency shift between magnon Stokes and anti-Stokes peaks  $\Delta f$  on magnon wavevector  $q$  allows us to determine IDMI constant  $D_s$  as:

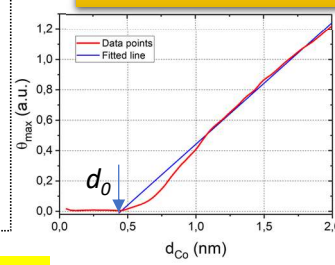
$$D_s = \frac{\pi M_s \Delta f}{2\gamma q} d_{co} \quad \text{where: } M_s \text{ is magnetization saturation, and } \gamma \text{ - gyromagnetic ratio.}$$

**Fig.1.** PMOKE remanence images of: a) sapphire/Pt/W/Co/Pt and b) sapphire/Pt/Co/W/Pt samples. Colour intensity is proportional to  $\theta_{REM}$ , defined in the Fig.2 inset. The regions with perpendicular magnetic anisotropy (PMA) are visible as light grey areas while the black regions illustrate either in-plane magnetization (thick  $d_{co}$ ) or nonmagnetic or superparamagnetic states (very thin  $d_{co}$ ). The blue and red dotted lines describe SRT positions of  $d_{SRT}$  and the thickness of magnetic dead layer  $d_0$  as a function of  $d_w$ . The anisotropy fields in Fig. 5 were determined along dashed horizontal lines (Fig.1a).

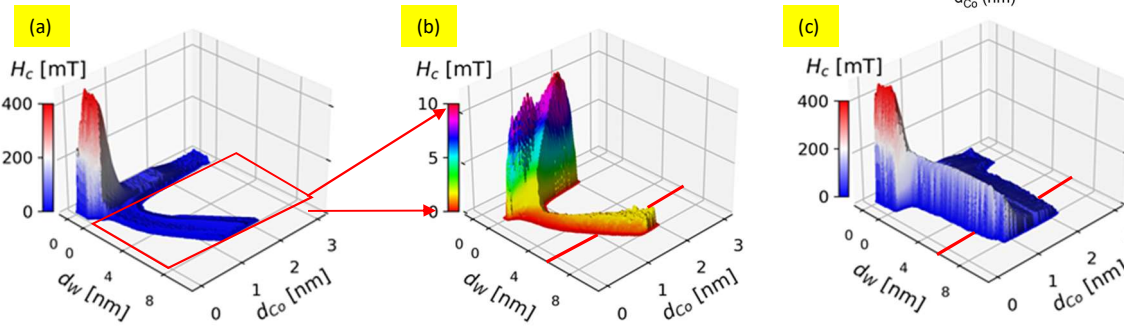


**Fig.2.** Exemplary PMOKE hysteresis loops recorded in the Pt/W/Co/Pt sample at the points  $\alpha$ ,  $\beta$ ,  $\gamma$  and  $\delta$  along the green vertical line ( $d_w=6.1$ nm, Fig. 1a). PMOKE rotation angle  $\theta$  is proportional to the local out-of-plane magnetization.  $\alpha$ )  $d_{co}=0.18$  nm - non-magnetic area;  $\beta$ )  $d_{co}=0.62$  nm - non-hysteresis (super)paramagnetic area;  $\gamma$ )  $d_{co}=1.17$  nm - "rectangular" shape loop typical for the PMA;  $\delta$ )  $d_{co}=1.87$  nm - non-hysteresis loop typical for magnetization in-plane. The definitions of  $\theta_{max}$ ,  $\theta_{SRT}$ , and  $H_c$  values, used in further calculations, are indicated by the red dots in the inset.

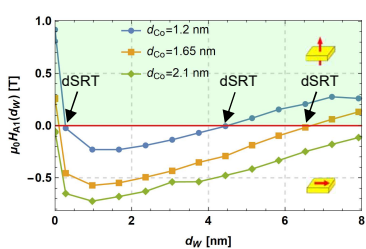
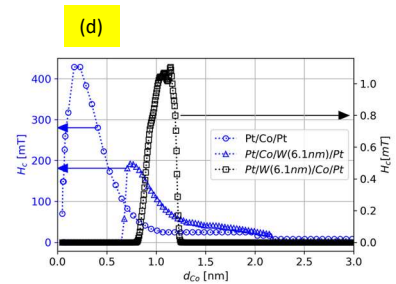
## Results & Discussions



**Fig.3.** An exemplary dependence of  $\theta_{max}(d_{co})$  for  $d_w=7.9$  nm (sample Pt/W/Co/Pt). The intersection of linear approximation (blue line) with  $\theta_{max}=0$  corresponds to the magnetic dead layer thickness  $d_0$ .



**Fig.4** 3D maps of coercive fields  $H_c(d_w, d_{co})$  in PMA regions. Figs (a) and (c) show the full ranges of  $H_c$  in Pt/W/Co/Pt and Pt/Co/W/Pt samples, respectively. High coercivity regions (red capped) are recorded in the reference Pt/Co/Pt regions. Fig (b) presents magnified part of (a). Fig (d) shows  $H_c$  a profile in the reference region and for a W buffer/cover layer thickness of  $d_w=6.1$  nm.



**Fig.5.** The dependences of anisotropy field  $H_{A1}(d_w)$  for  $d_{co}=1.2, 1.65$  and  $2.1$  nm (according to the profile lines shown in Fig.1A with appropriate colors).

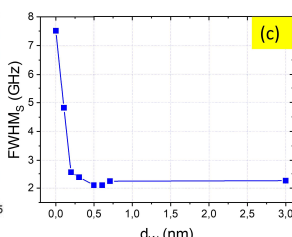
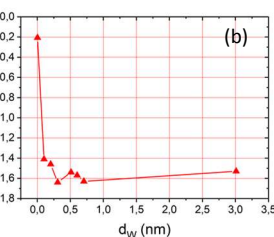
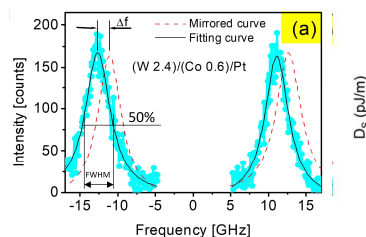
## Conclusions:

Abrupt changes (induced by sub-/few W atomic monolayer) of :

- magnetic anisotropy field  $H_{A1}$
- coercivity field  $H_c$  (two orders or two time decrease in the case of W buffer and overlayer, respectively)
- decrease in IDMI constant  $D_s$  and width of Stokes peaks ( $FWHM_s$ )

The difference in the W buffer and overlayer influence on magnetic properties can be explained by modifications of both Co interfaces and crystallographic structure.

W layer-tuned properties are important for designing more complicated nanostructures (synthetic ferromagnets and antiferromagnets with IDMI) which could be perspective for further applications, e.g. magnonic structures.



References :

1. A. Fert, et al., Nat. Rev. Mater. **2** (7), 17031 (2017)
2. W. Legrand, et al., Phys. Rev. Materials **6**, 024408 (2022)
3. Z. Kurant, et al., J. Magn. Magn. Mater. **558**, 169485 (2022)



## Acknowledgements:

This work is supported by the National Science Center in Poland under the project OPUS-19 (2020/37/B/ST/5/02299).

email: wawro@ifpan.edu.pl