

Pair-superfluid phase of dipolar molecules in optical lattice

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Abstract

We study the extended Bose–Hubbard model describing an ultra-cold gas of dipolar molecules in an optical lattice, taking into account all on-site and nearest-neighbor interactions, including occupation-dependent tunneling and pair tunneling terms. We show that these terms lead to additional quantum phase transitions and can destroy insulating phases. These considerable changes of the phase diagram have to be taken into account in upcoming experiments with dipolar molecules.

Bosons in optical lattice 1

• Hamiltonian of the system

$$\mathcal{H} = \int d^3r \Psi^\dagger(\mathbf{r}) \left[-\frac{\hbar^2}{2m} \nabla^2 + V_{\text{ext}}(\mathbf{r}) \right] \Psi(\mathbf{r}) + \frac{1}{2} \int \int d^3r d^3r' \Psi^\dagger(\mathbf{r}) \Psi^\dagger(\mathbf{r}') \mathcal{V}(\mathbf{r} - \mathbf{r}') \Psi(\mathbf{r}') \Psi(\mathbf{r})$$

• Lattice potential

$$V_{\text{ext}}(\mathbf{r}) = V_0 \left[\sin^2 \frac{2\pi x}{\lambda} + \sin^2 \frac{2\pi y}{\lambda} \right] + \frac{m\Omega^2}{2} z^2$$

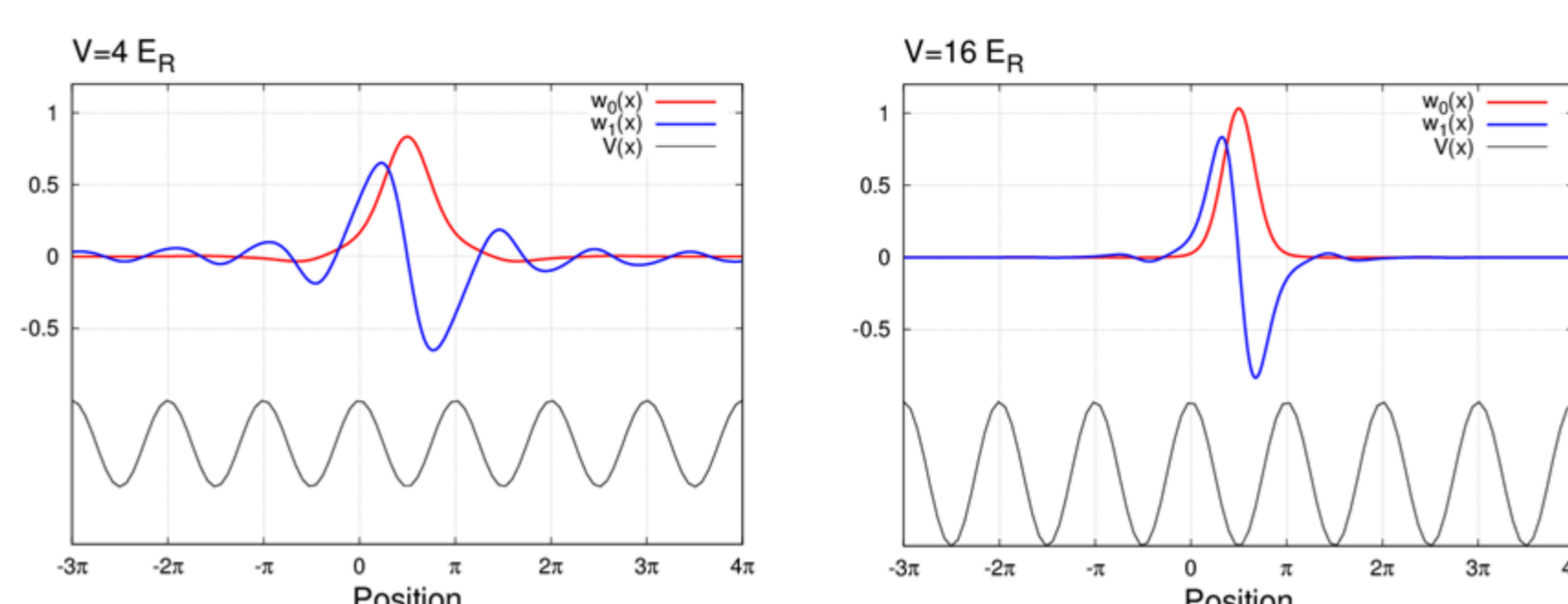
• Natural units of the problem

- laser wave length $\lambda/2\pi$ - lattice flattening $\kappa = \frac{\hbar\Omega}{2E_R}$
- recoil energy $E_R = \frac{2\pi^2\hbar^2}{m\lambda^2}$

Wannier functions 2

• Convenient basis

$\mathcal{W}_i^\alpha(\mathbf{r})$ wave function localized in i -th lattice site
 α denotes appropriate Bloch band



Field operator decomposition 3

$$\Psi(\mathbf{r}) = \sum_\alpha \sum_i \hat{a}_i^{(\alpha)} \mathcal{W}_i^\alpha(\mathbf{r}) \approx \sum_i \hat{a}_i \mathcal{W}_i(\mathbf{r})$$

the lowest band approximation

• Single particle Hamiltonian

$$\mathcal{H} = \int d^3r \Psi^\dagger(\mathbf{r}) \left[-\frac{\hbar^2}{2m} \nabla^2 + V_{\text{ext}}(\mathbf{r}) \right] \Psi(\mathbf{r}) = \sum_{ij} \hat{a}_i^\dagger \hat{a}_j \int d^3r \mathcal{W}_i^*(\mathbf{r}) \left[-\frac{\hbar^2}{2m} \nabla^2 + V_{\text{ext}}(\mathbf{r}) \right] \mathcal{W}_j(\mathbf{r}) = \mathbf{E} \sum_i \hat{a}_i^\dagger \hat{a}_i - \mathbf{J} \sum_i \sum_{\langle j \rangle} \hat{a}_i^\dagger \hat{a}_j + \dots$$

summation over nearest neighbours of i

Long range interactions 4

$$\Psi(\mathbf{r}) = \sum_\alpha \sum_i \hat{a}_i^{(\alpha)} \mathcal{W}_i^\alpha(\mathbf{r}) \approx \sum_i \hat{a}_i \mathcal{W}_i(\mathbf{r})$$

lowest band approximation

• Interaction Hamiltonian

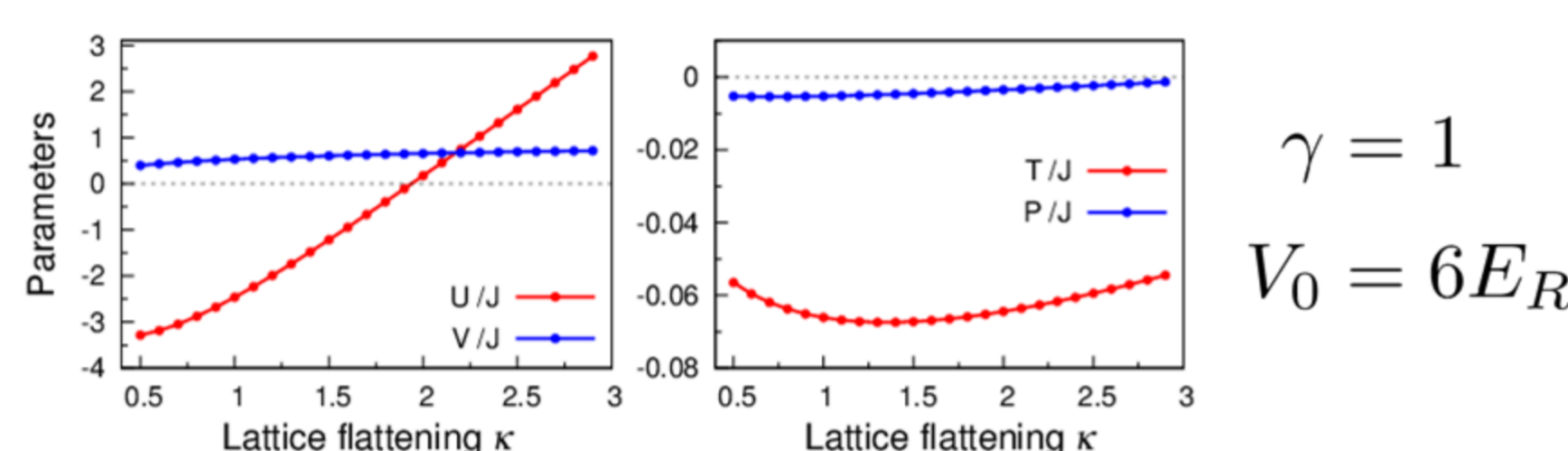
$$\mathcal{H} = \frac{1}{2} \int \int d^3r d^3r' \Psi^\dagger(\mathbf{r}) \Psi^\dagger(\mathbf{r}') \mathcal{V}(\mathbf{r} - \mathbf{r}') \Psi(\mathbf{r}') \Psi(\mathbf{r}) = \frac{1}{2} \sum_{ijkl} \hat{a}_i^\dagger \hat{a}_j^\dagger \hat{a}_k \hat{a}_l \int \int d^3r d^3r' \mathcal{W}_i^*(\mathbf{r}) \mathcal{W}_j^*(\mathbf{r}') \mathcal{V}(\mathbf{r} - \mathbf{r}') \mathcal{W}_k(\mathbf{r}') \mathcal{W}_l(\mathbf{r}) = \frac{\mathbf{U}}{2} \sum_i \hat{a}_i^\dagger \hat{a}_i^\dagger \hat{a}_i \hat{a}_i + \mathbf{V} \sum_i \sum_{\langle j \rangle} \hat{a}_i^\dagger \hat{a}_j^\dagger \hat{a}_i \hat{a}_j + \dots - \mathbf{T} \sum_i \sum_{\langle j \rangle} \hat{a}_i^\dagger (\hat{a}_i^\dagger \hat{a}_i + \hat{a}_j^\dagger \hat{a}_j) \hat{a}_j + \frac{\mathbf{P}}{2} \sum_i \sum_{\langle j \rangle} \hat{a}_i^\dagger \hat{a}_i^\dagger \hat{a}_j \hat{a}_j + \dots$$

Studied Hamiltonian 5

$$\mathcal{H} = \mathbf{E} \sum_i \hat{a}_i^\dagger \hat{a}_i - \mathbf{J} \sum_{\langle ij \rangle} \hat{a}_i^\dagger \hat{a}_j + \frac{\mathbf{U}}{2} \sum_i \hat{n}_i (\hat{n}_i - 1) + \mathbf{V} \sum_{\langle ij \rangle} \hat{n}_i \hat{n}_j - \mathbf{T} \sum_{\langle ij \rangle} \hat{a}_i^\dagger (\hat{n}_i + \hat{n}_j) \hat{a}_j + \frac{\mathbf{P}}{2} \sum_{\langle ij \rangle} \hat{a}_i^\dagger \hat{a}_i^\dagger \hat{a}_j \hat{a}_j$$

• Considered interaction

$$\mathcal{V}(\mathbf{r} - \mathbf{r}') = \gamma \left(\frac{1}{|\mathbf{r} - \mathbf{r}'|^3} - 3 \frac{(z - z')^2}{|\mathbf{r} - \mathbf{r}'|^5} \right) + g \delta^{(3)}(\mathbf{r} - \mathbf{r}')$$



Physical realization 6

• Two dimensionless parameters

$$g = 16\pi^2 \frac{\alpha_s}{\lambda} \quad \gamma = \frac{d^2}{\hbar^2 \epsilon_0 \lambda}$$

s-wave scattering length electric dipole moment

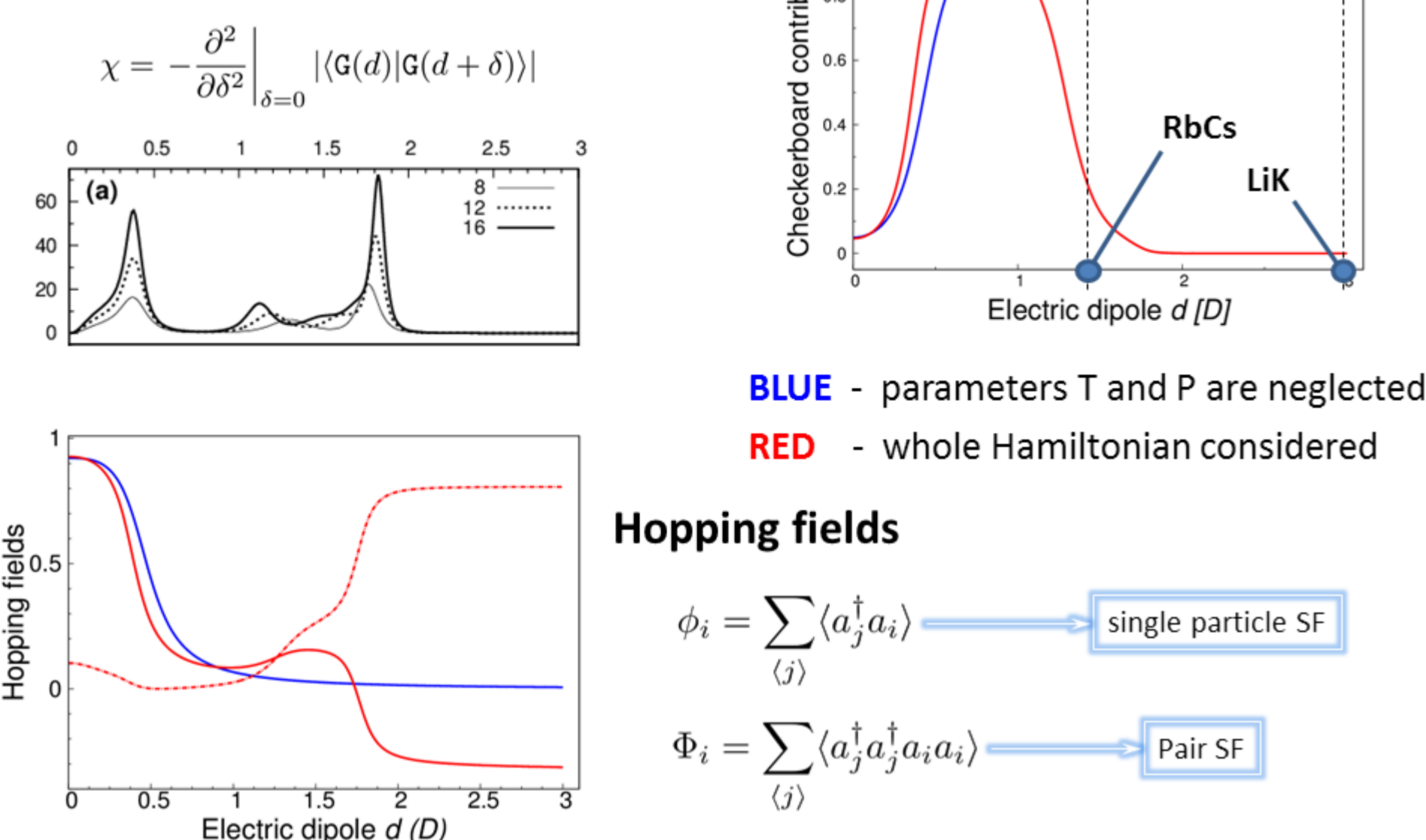
• Model assumptions

- optical lattice laser wave length $\lambda = 790$ nm
- mass \sim mass of the RbCs molecule
- scattering length ~ 100 Bohr Radius $\longrightarrow g \approx 1$
- dipole moment up to 3 debye $\longrightarrow \gamma$ up to 470

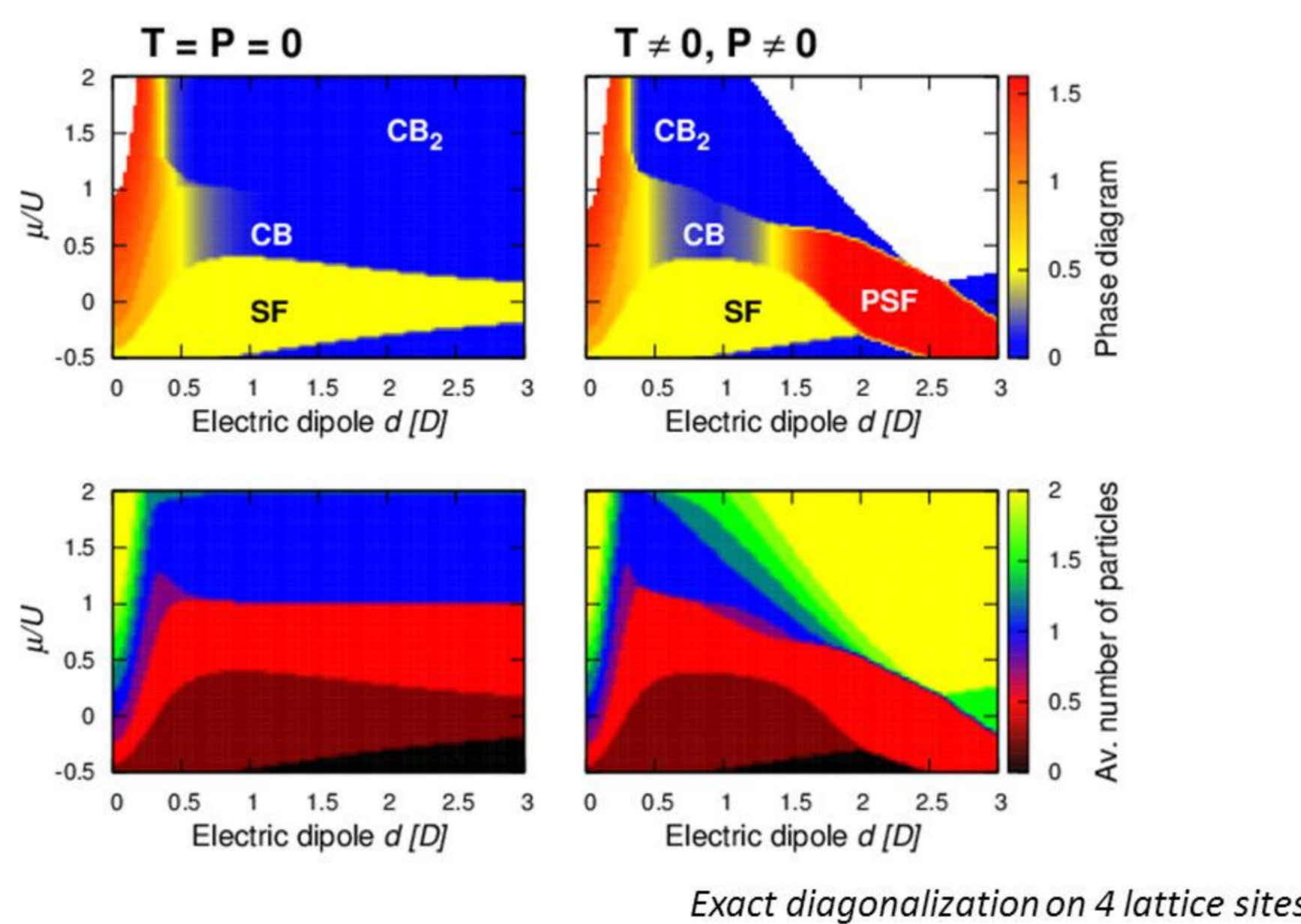
Half filled 1D system 7

METHOD:
Exact diagonalization of the 1D Hamiltonian with $N = 8, 12, 16$ sites

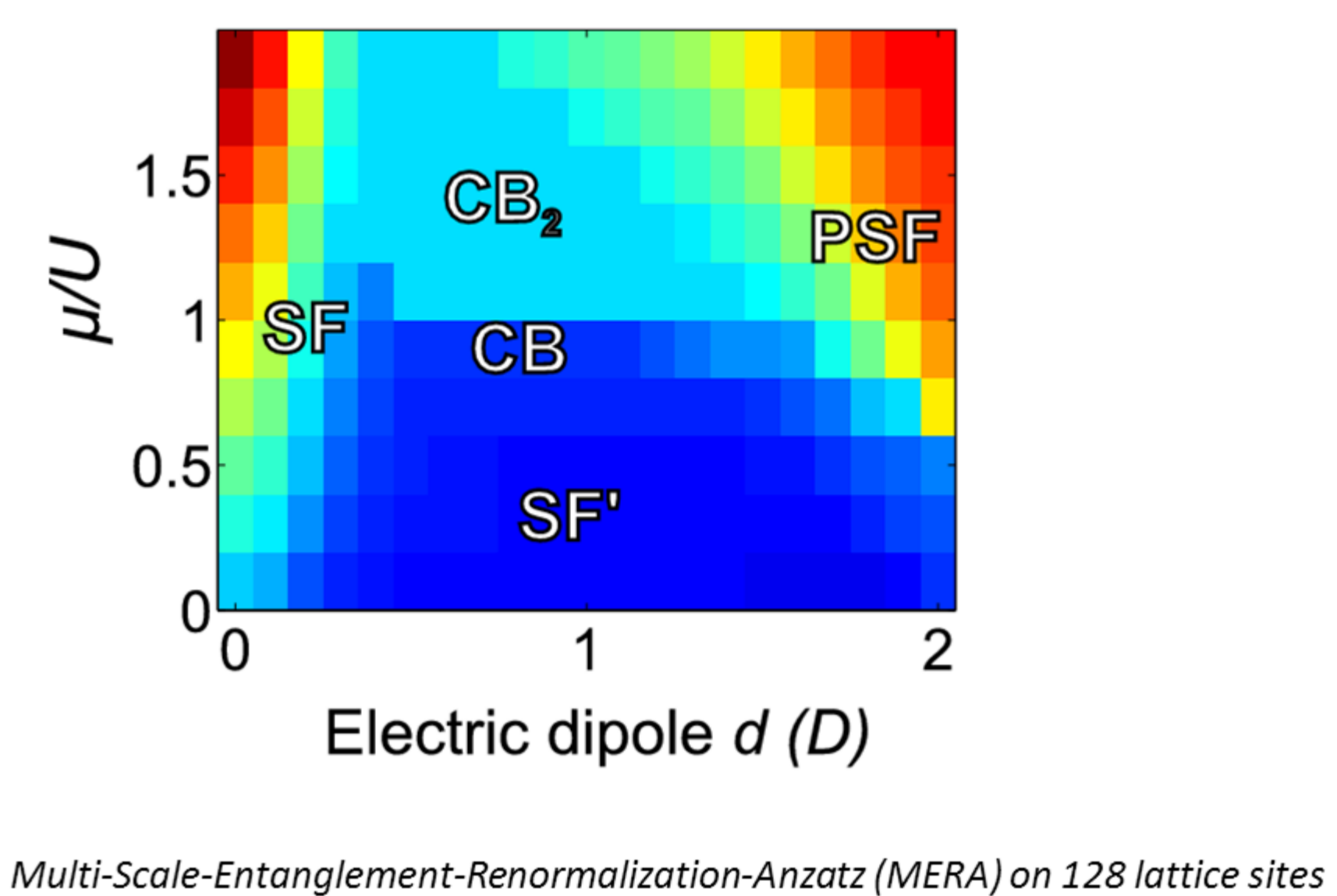
• Susceptibility



Grand canonical phase diagram 8



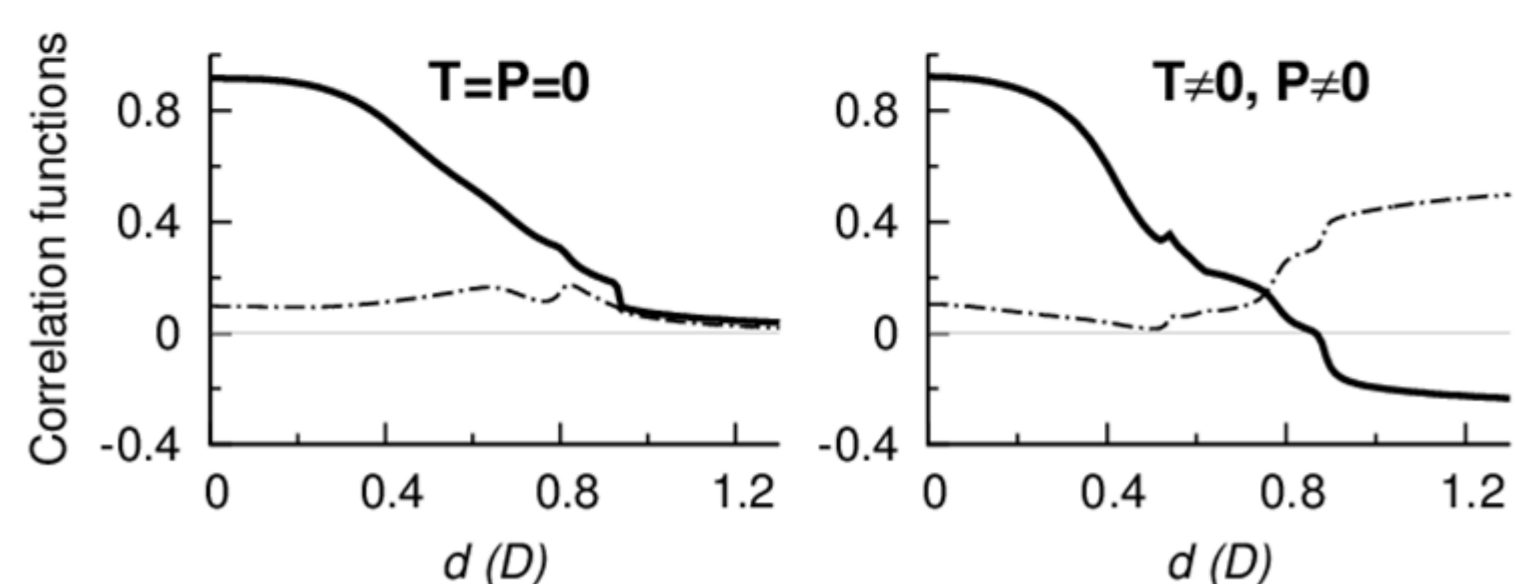
Grand canonical phase diagram 9



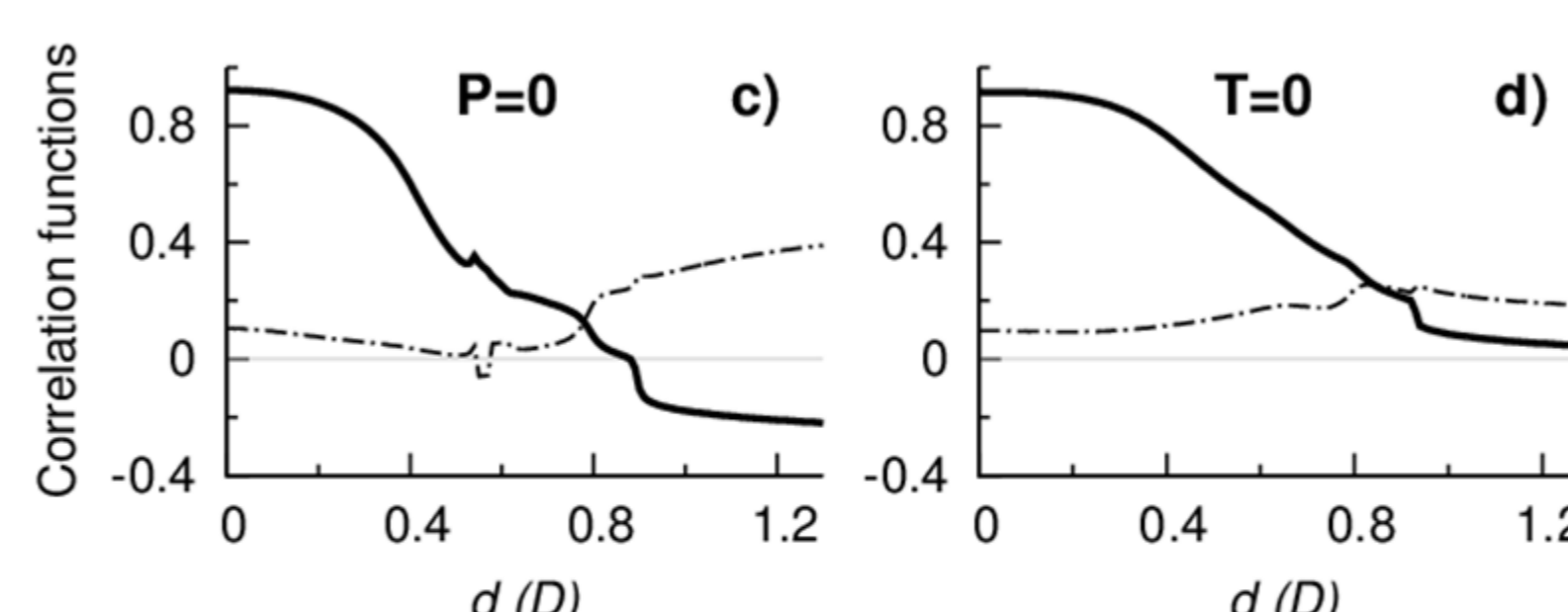
Further neighbours 10

• generalization of the Hamiltonian

$$\mathbf{V} \sum_{\langle ij \rangle} \hat{n}_i \hat{n}_j \longrightarrow \sum_{L=1}^{N/2-1} \mathbf{V}_L \sum_i \hat{n}_i \hat{n}_{i+L}$$



Origins of pair-superfluidity 11



We have further checked that counter-intuitively PSF arises predominantly due to correlated tunneling T. Without this term PSF phase can not be reached for reasonable electric moments.

Summary

Commonly neglected terms in the extended Bose-Hubbard model for dipolar molecules can lead to NEW PHENOMENA

MORE DETAILS:

T. Sowiński, O. Dutta, P. Hauke, L. Tagliacozzo, M. Lewenstein
"Dipolar molecules in optical lattices"
Physical Review Letters (in press)