

**Conclusions:** a) our experiment saw survival of the expanded depletion apparently without need for impurities  
 b) the apparent contact (tail strength) in the experiment was again enhanced compared to in situ values  
 c) survival and contact strengthening can also be seen in simulations, but less strongly  
 d) survival relies on an asymmetric trap and fast, adiabatic release, mostly from the edges of the BEC  
 e) tail strength appears to increase due to acceleration by the mean field during escape from the BEC

## Quantum depletion and Tan's contact

### Uncondensed atoms in a BEC

Consider a uniform piece of the gas

$$\hat{b}_k^\dagger = u_k \hat{a}_k^\dagger + v_k \hat{a}_{-k}$$

$$u_k = \cosh \theta_k, \quad v_k = \sinh \theta_k$$

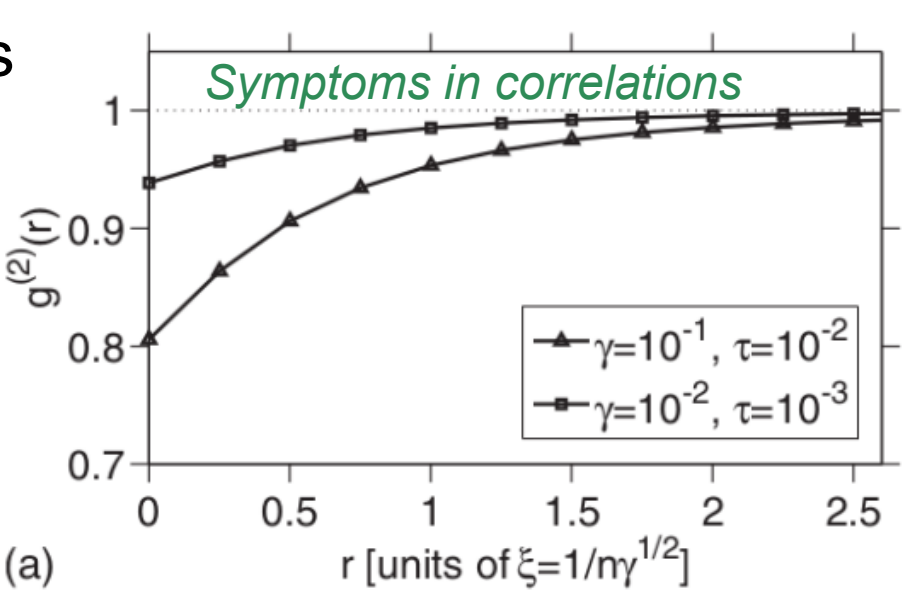
$$\theta_k = \frac{1}{2} \log \frac{\hbar^2 k^2 / 2m}{\epsilon(k)} < 0$$

$$\epsilon(k) = \sqrt{\left(\frac{\hbar^2 k^2}{2m}\right)^2 + gn \frac{\hbar^2 k^2}{m}}$$

$$\rho(\mathbf{k}) = \langle \hat{a}_k^\dagger \hat{a}_k \rangle \quad \text{depletion}$$

$$= (u_k^2 + v_k^2) \langle \hat{b}_k^\dagger \hat{b}_k \rangle + v_k^2$$

$$\langle \hat{b}_k^\dagger \hat{b}_k \rangle = (\exp[\epsilon(k)/k_B T] - 1)^{-1}$$



$$\lim_{k \rightarrow \infty} \rho(\mathbf{k}) \propto k^{-4}$$

Tan's contact  
 $C = \lim_{k \rightarrow \infty} k^4 \rho(k)$

In the non-interacting ( $a \rightarrow 0$ ) limit,  $u_k = 1$  and  $v_k = 0$ ,

$$\text{in situ tails: } \lim_{k \rightarrow \infty} n(k) = \frac{C}{k^4} = \frac{64\pi^2 a^2 N_0 n_0}{7 k^4}$$

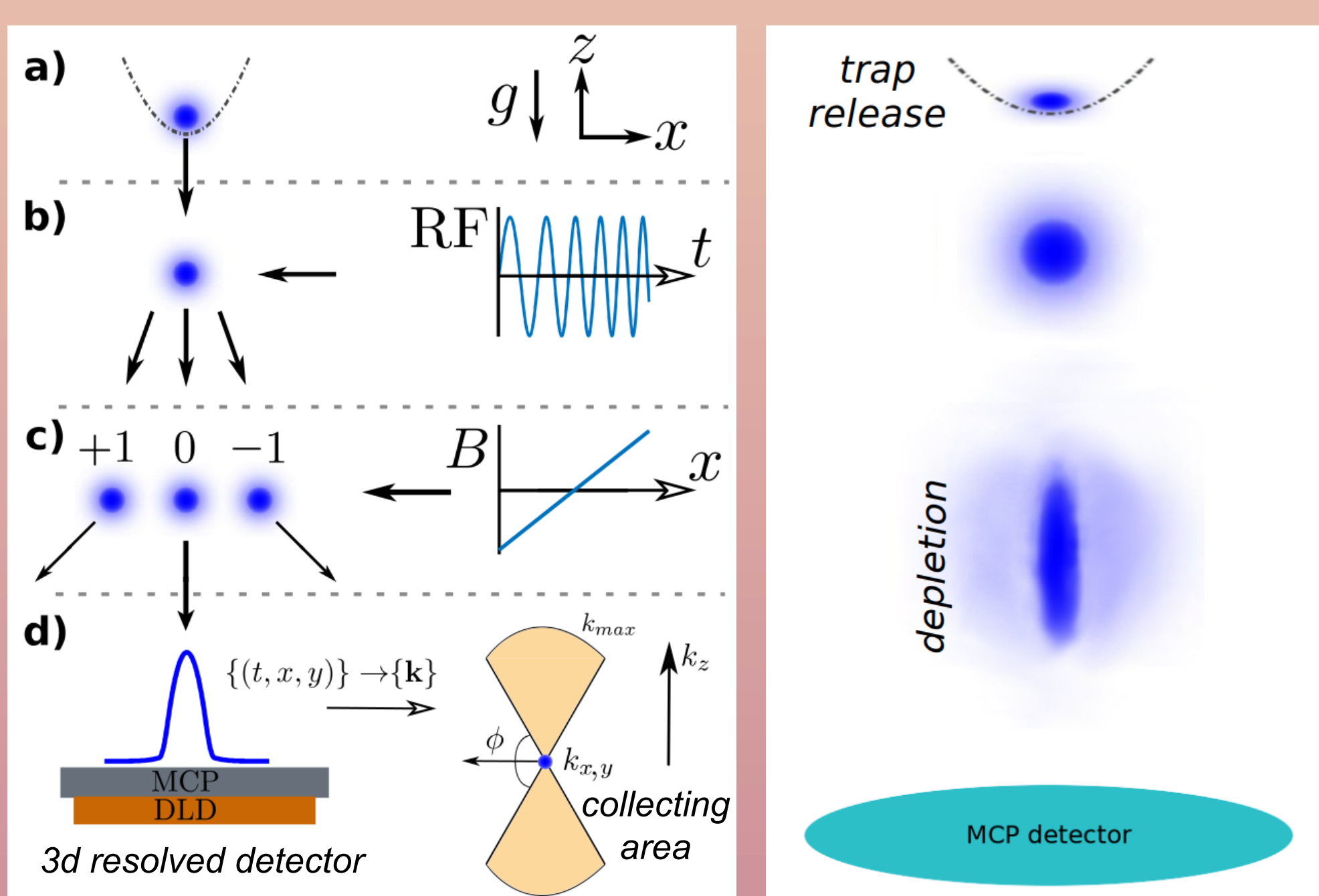
## The mystery

- Palaiseau experiments saw stronger tails than expected *in situ*. (David Clement's group, PRL 2016)
- Theory suggests no survival possible (Qu, Pitaevskii, Stringari, PRA 94, 063635 (2016))
- Later Palaiseau experiments indicate survival depends on presence of impurities (Cayla et al arXiv:2204.10697, David Clement's group)

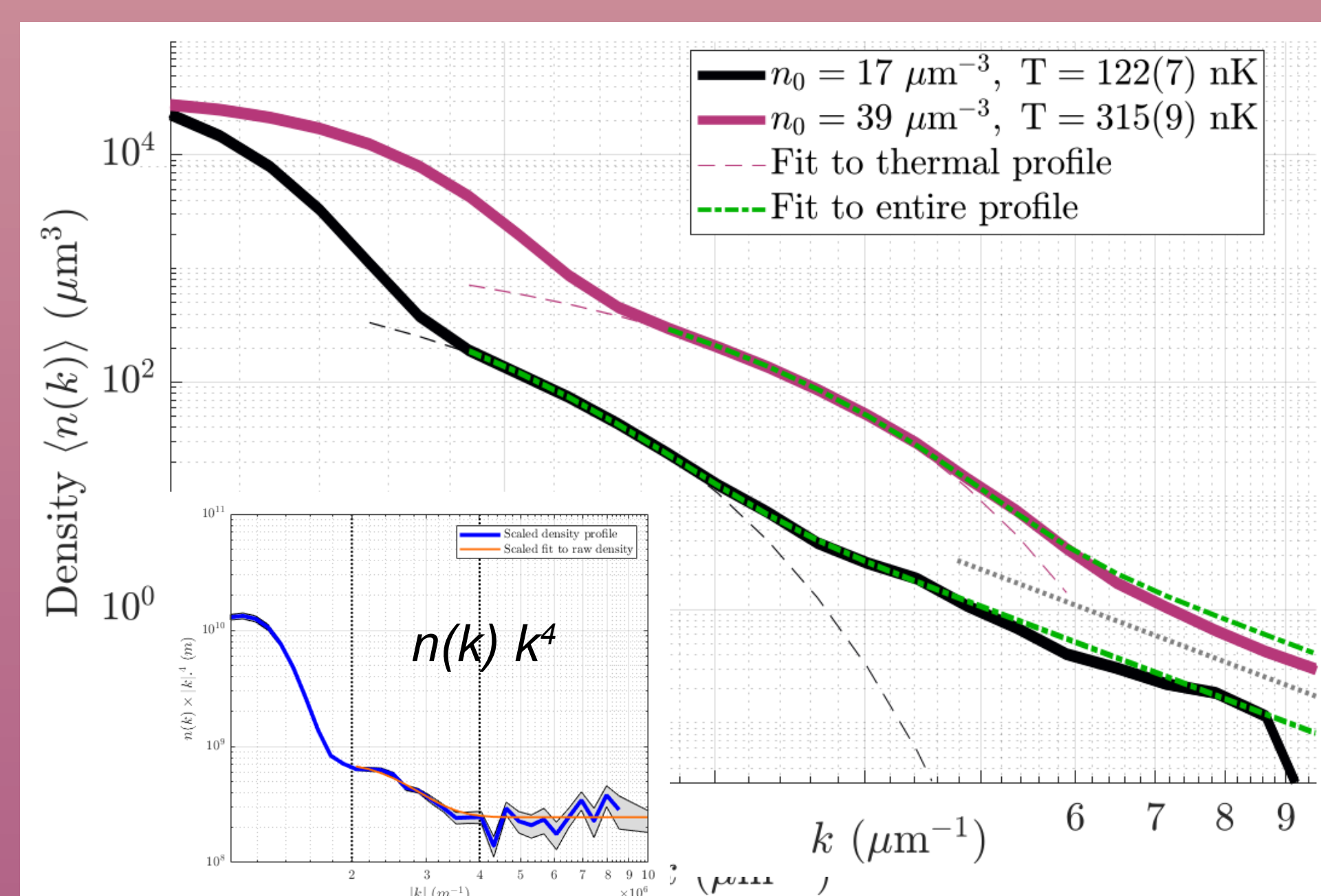
## Our aims

- Canberra expt. has a fully magnetic trap → only one component trapped → lack of impurities *in situ*. What about tails?
- Theory: can we simulate time evolution and see how survival is possible?

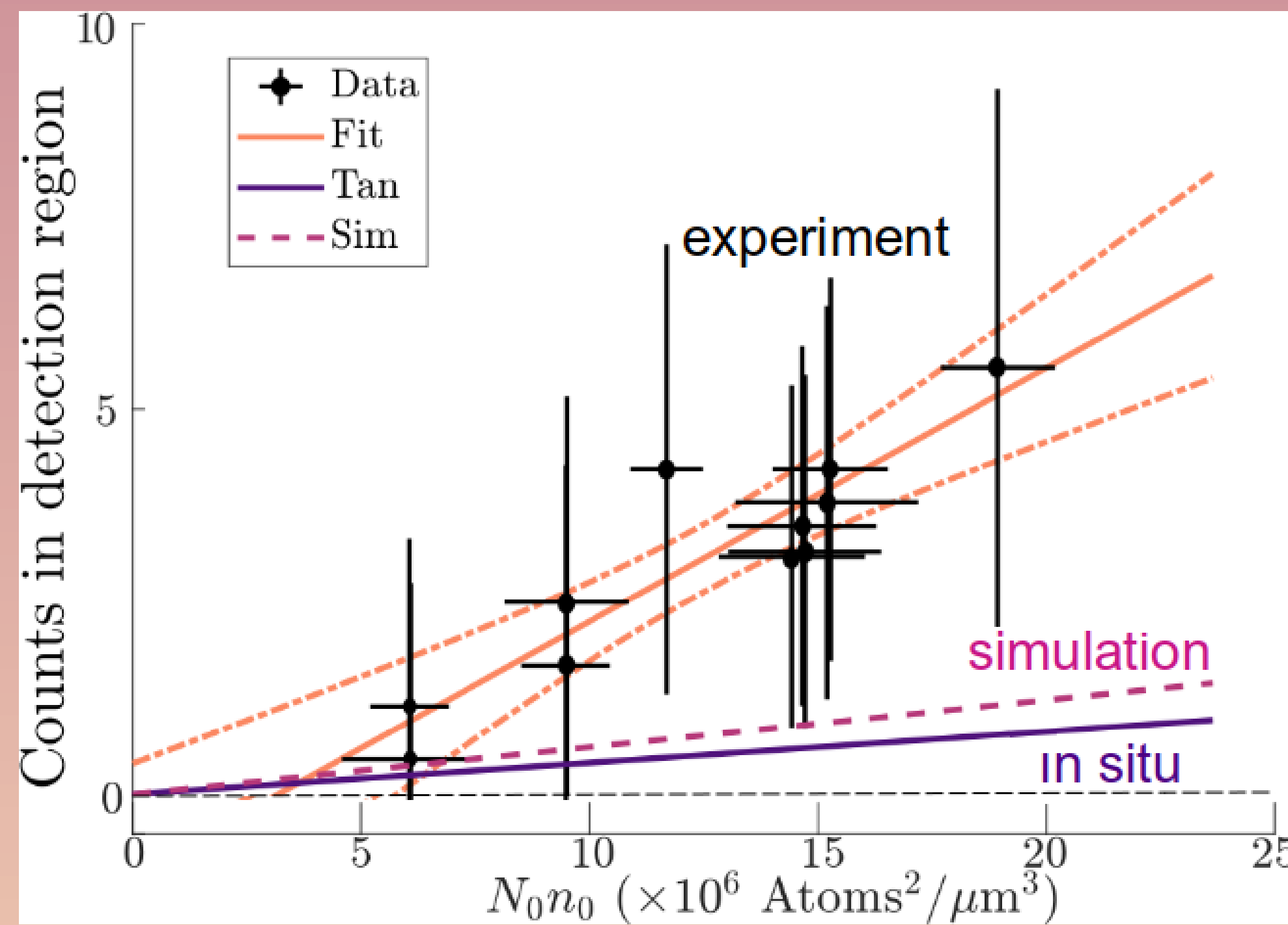
## Experiment



## Velocity profile on MCP detector



## Experimental tail strength



## Simulation method

$$\hat{\Psi}(\mathbf{x}, t) = \phi(\mathbf{x}, t) + \hat{\Psi}_B(\mathbf{x}, t) \quad \text{Bogoliubov}$$

PD, Chwedeńczuk, Ziń, Trippenbach, PRA 83, 063625 (2011)

plain "STAB" was useful for BEC collisions when Bogoliubov modes were separated from the condensate in k-space  
 Now this is no longer true (lots of overlap)

However, need to also impose orthogonality this time

$$\int d^3x \hat{\Psi}_B^\dagger(\mathbf{x}, t) \phi(\mathbf{x}, t) = 0$$

Proper STAB equations that preserve orthogonality

$$i\hbar \frac{d\phi}{dt} = \left[ -\frac{\hbar^2}{2m} \nabla^2 + g|\phi|^2 + V(\mathbf{x}, t) \right] \phi \quad \text{GPE mean field } \mathcal{H}(g, \phi) \psi_B$$

$$i\hbar \frac{d\psi_B}{dt} = \mathcal{H}(g, \phi) \psi_B + \mathcal{P}_\perp \left\{ g_B |\phi|^2 \psi_B + g_B \phi^2 \tilde{\psi}_B^* + \sqrt{-ig_B} \phi \xi(\mathbf{x}, t) \right\}$$

$$i\hbar \frac{d\tilde{\psi}_B}{dt} = \mathcal{H}(g, \phi) \tilde{\psi}_B + \mathcal{P}_\perp \left\{ g_B |\phi|^2 \tilde{\psi}_B + g_B \phi^2 \psi_B^* + \sqrt{-ig_B} \phi \tilde{\xi}(\mathbf{x}, t) \right\}$$

Thankfully, projection can be done very efficiently

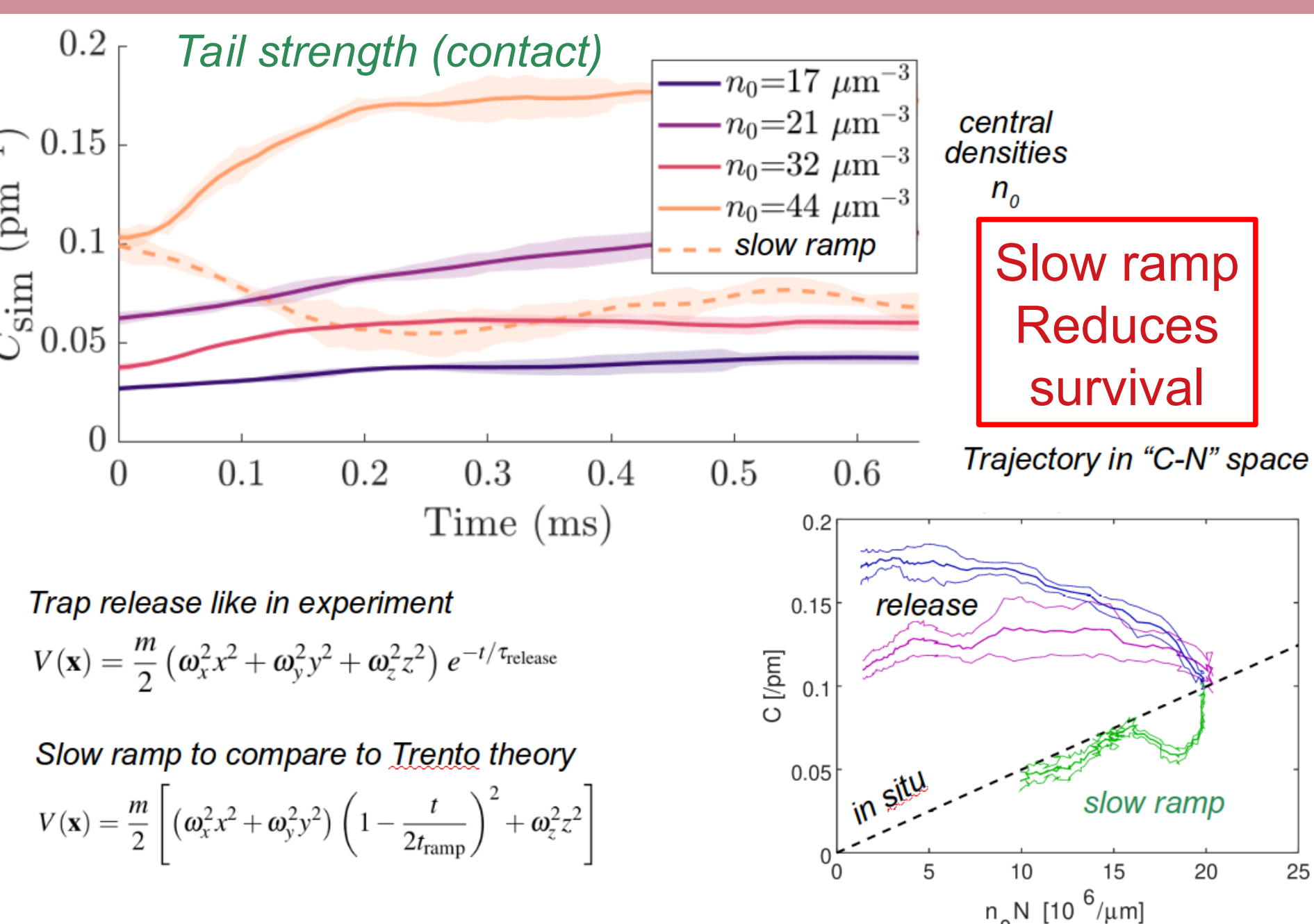
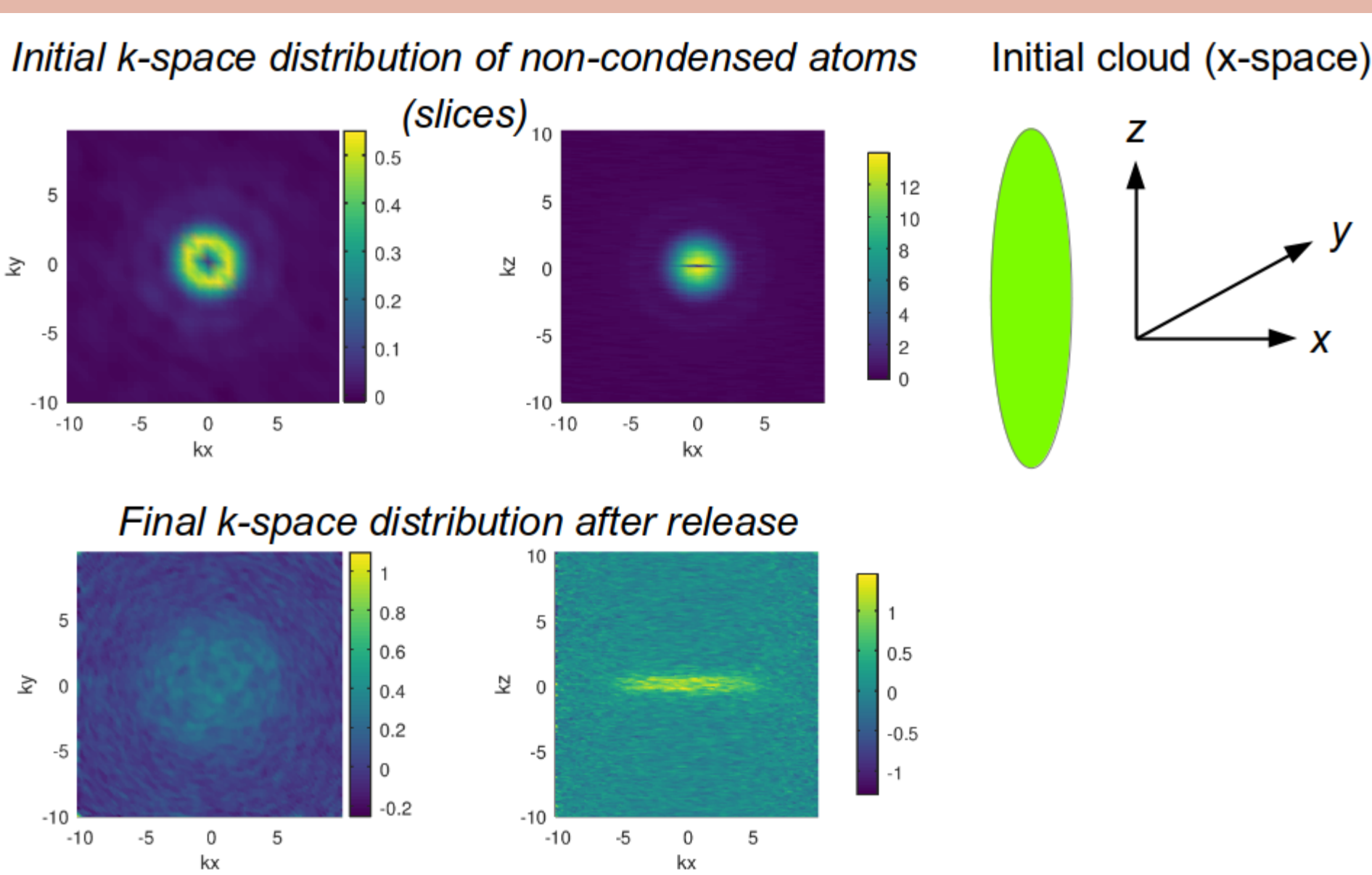
$$\mathcal{P}_\perp f(\mathbf{x}) = f(\mathbf{x}) - \frac{1}{N} \left[ \int d^3x' \phi(\mathbf{x}')^* f(\mathbf{x}') \right] \phi(\mathbf{x})$$

$$\text{Gaussian real white noise } \langle \xi(\mathbf{x}, t) \xi(\mathbf{y}, t') \rangle = \delta^3(\mathbf{x} - \mathbf{y}) \delta(t - t')$$

Depletion density

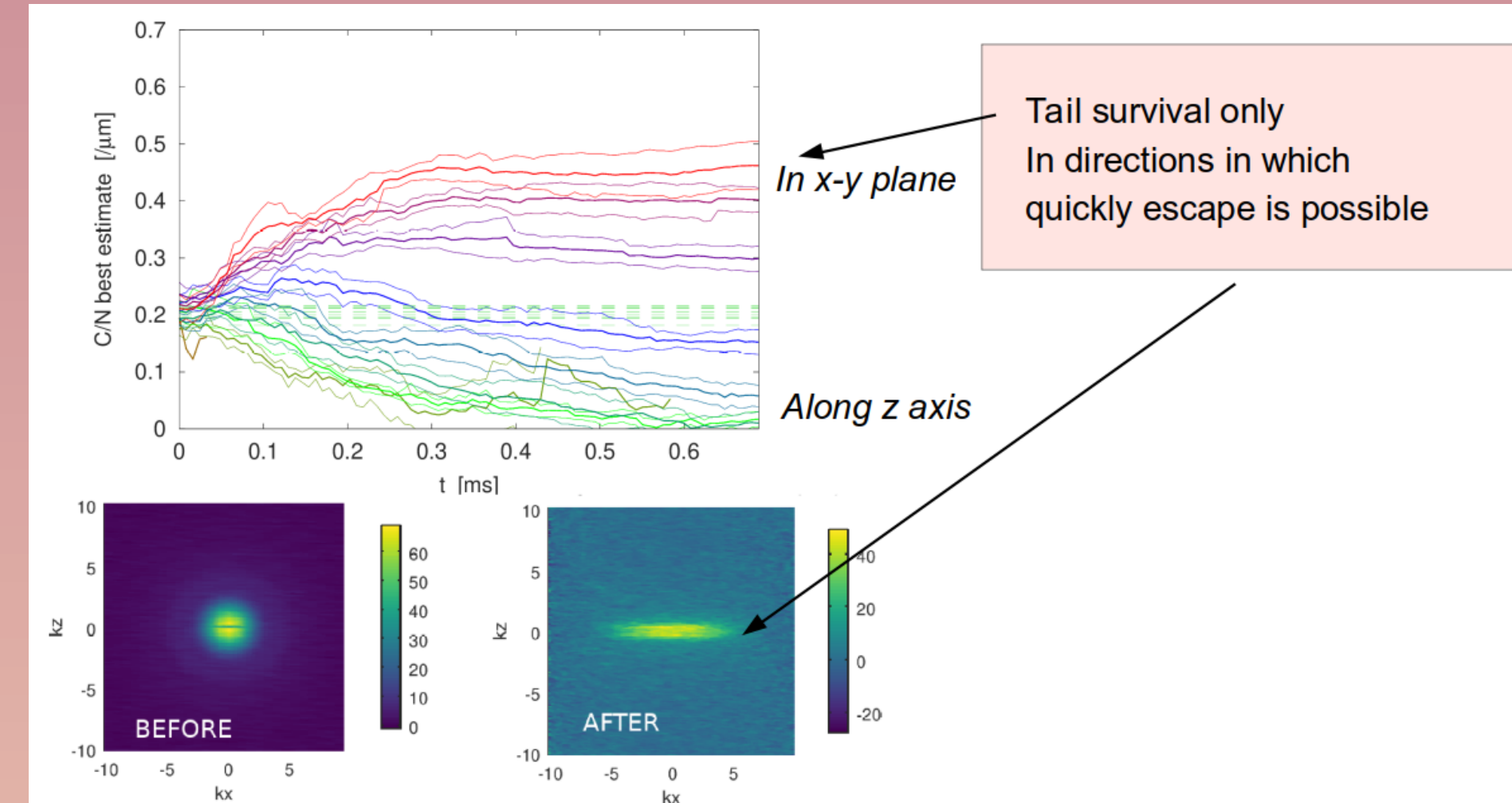
$$n_B(\mathbf{k}) = \text{Re} \langle \tilde{\psi}_B^\dagger(\mathbf{k}, t) \psi_B(\mathbf{k}, t) \rangle_{\text{stoch}}$$

## Simulation: distribution, time evolution

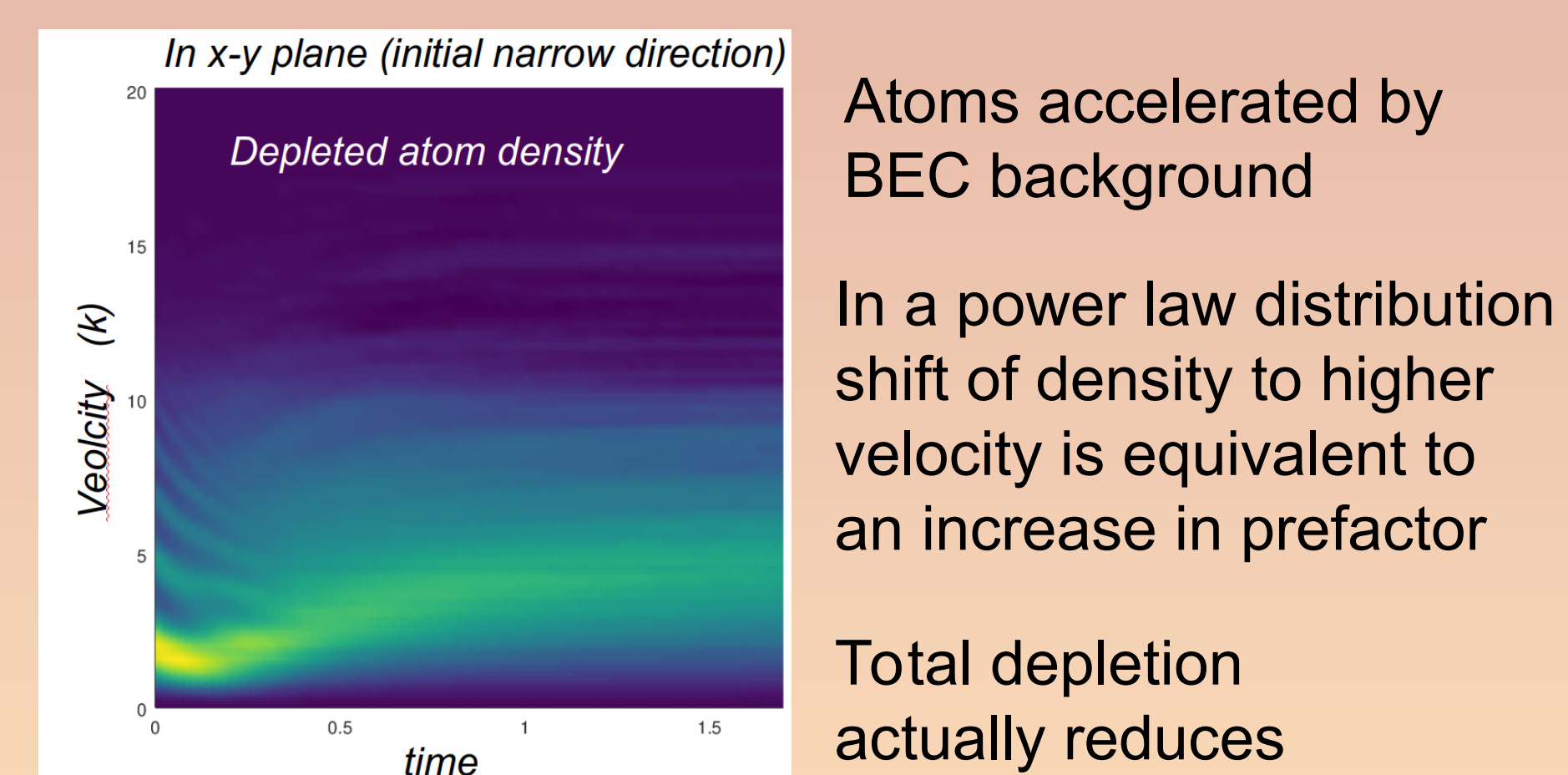


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## Anisotropy - simulation



## Apparent growth of tails



## Toy model - survival process

Single pair of modes in local density  $n(t)$

$$\frac{d}{dt} \begin{bmatrix} \hat{a}_k(t) \\ \hat{a}_{-k}^\dagger(t) \end{bmatrix} = -\frac{i}{\hbar} \begin{bmatrix} \hbar^2 k^2 / 2m + 2gn(t) - \mu(t) & gn(t) \\ -gn(t) & -\hbar^2 k^2 / 2m - 2gn(t) + \mu(t) \end{bmatrix} \begin{bmatrix} \hat{a}_k(t) \\ \hat{a}_{-k}^\dagger(t) \end{bmatrix}$$

Initial radial position  
 $R_0 = \frac{\sqrt{y^2 + z^2}}{R_\perp} = \frac{\omega_\perp \sqrt{y^2 + z^2}}{\sqrt{2gn_0/m}}$

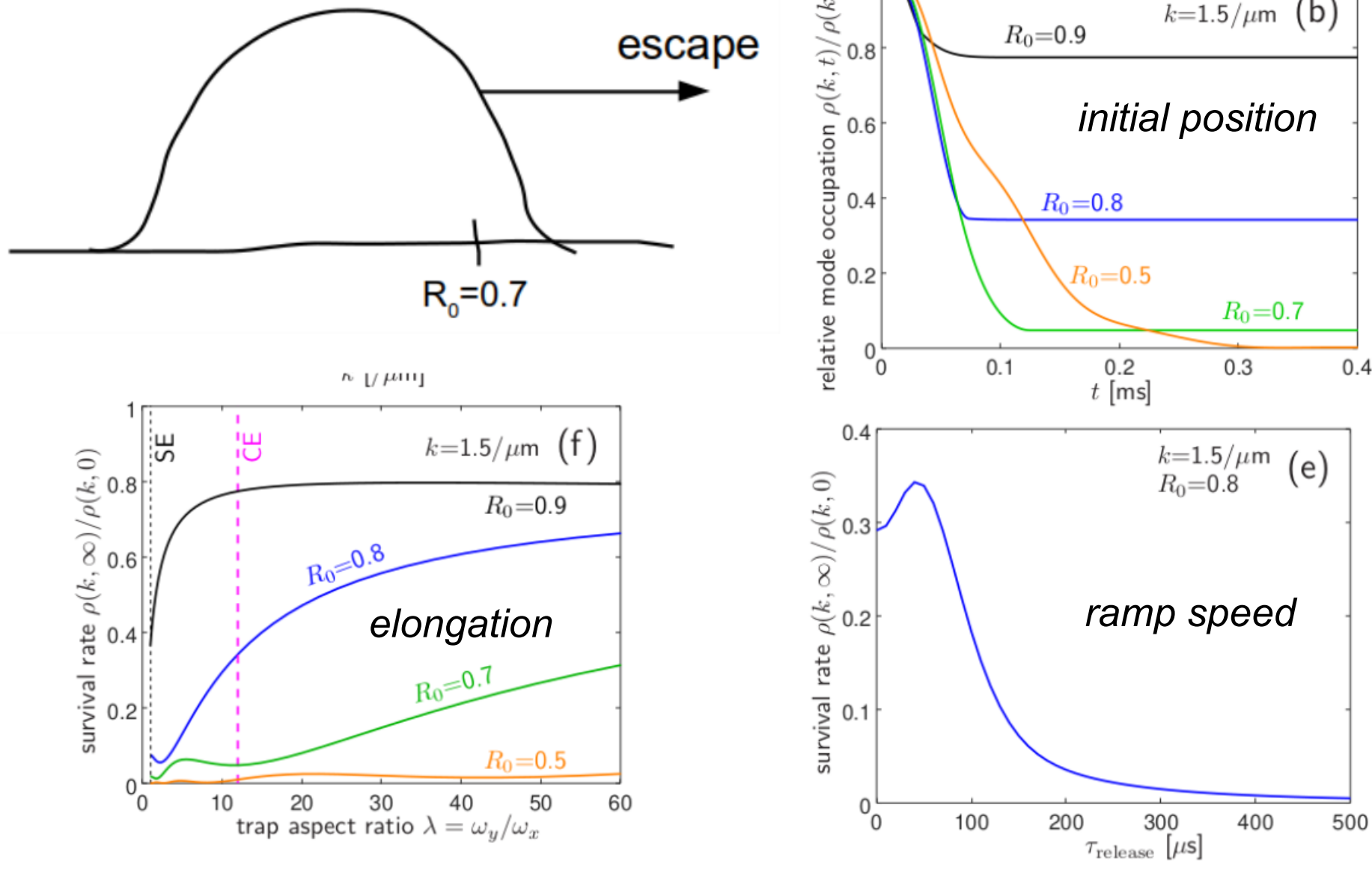
Estimated radial position at time  $t$  (flight + riding condensate expansion)  
 $r(t) = R_0 R_\perp \sqrt{1 + \omega_\perp^2 t^2} + \frac{\hbar k t}{m}$

Radial position relative to condensate  
 $r_{\text{rel}}(t) = \frac{r(t)}{\sqrt{1 + \omega_\perp^2 t^2}}$

Working backwards to initial radial position estimate  
 $n_c(r, t) = \frac{n_c(r, 0)}{(1 + \omega_\perp^2 t^2) \sqrt{1 + \omega_\perp^2 t^2}}$ , where  $n_c(r, 0) = \begin{cases} n_0 [1 - (r/R_\perp)^2] & \text{if } r < R_\perp \\ 0 & \text{if } r \geq R_\perp \end{cases}$

Ramped BEC density  
 $n(t) = n(0) e^{-3t/\tau_{\text{release}}} + (1 - e^{-3t/\tau_{\text{release}}}) n_c(r, 0, t)$

## Plots show survival rates



## Final comparison of tail strengths

	tail strength vs value <i>in situ</i>	
Chang, Bouton, Cayla, Qu, Aspect, Westbrook, Clement, PRL 117, 235303 (2016)	$\times 6 \pm 1$	Experiment (Palaiseau)
Qu, Pitaevskii, Stringari PRA 94, 063635 (2016)	0	Theory (Trento)
Ross, Deuar, Shin, Thomas, Henson, Hodgman, Truscott, Sci Rep 12, 13178 (2022)	$\times 2 \pm 0.2$	Simulations (Warsaw)
Ross, Deuar, Shin, Thomas, Henson, Hodgman, Truscott, Sci Rep 12, 13178 (2022)	$\times 5 \pm 3$	Experiment (Canberra)
Cayla, Massignan, Giamarchi, Aspect, Westbrook, Clement to appear (2022)	$\times 0 - 6$ depending on impurities	Experiment 2 (Palaiseau)